

Volume Rendering Digest (for NeRF)

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Neural Radiance Fields [3] employ simple volume rendering as a way to overcome the challenges of differentiating through ray-triangle intersections by leveraging a probabilistic notion of visibility. This is achieved by assuming the scene is composed by a cloud of light-emitting particles whose density changes in space (in the terminology of physically-based rendering, this would be described as a volume with absorption and emission but no scattering [4, Sec 11.1]). In what follows, for the sake of exposition simplicity, and without loss of generality, we assume the emitted light *does not* change as a function of view-direction. This technical report is a condensed version of previous reports [1, 2], but rewritten in the context of NeRF, and adopting its commonly used notation¹.

Transmittance. Let the density field $\sigma(\mathbf{x})$, with $\mathbf{x} \in \mathbb{R}^3$ indicate the differential likelihood of a ray hitting a particle (i.e. the probability of hitting a particle while travelling an infinitesimal distance). We reparameterize the density along a given ray $\mathbf{r} = (\mathbf{o}, \mathbf{d})$ as a scalar function $\sigma(t)$, since any point \mathbf{x} along the ray can be written as $\mathbf{r}(t) = \mathbf{o} + t\mathbf{d}$. Density is closely tied to the transmittance function $\mathcal{T}(t)$, which indicates the probability of a ray traveling over the interval $[0, t)$ without hitting any particles. Then the probability $\mathcal{T}(t+dt)$ of not hitting a particle when taking a differential step dt is equal to $\mathcal{T}(t)$, the likelihood of the ray reaching t , times $(1 - dt \cdot \sigma(t))$, the probability of not hitting anything during the step:

$$\mathcal{T}(t + dt) = \mathcal{T}(t) \cdot (1 - dt \cdot \sigma(t)) \quad (1)$$

$$\frac{\mathcal{T}(t + dt) - \mathcal{T}(t)}{dt} \equiv \mathcal{T}'(t) = -\mathcal{T}(t) \cdot \sigma(t) \quad (2)$$

This is a classical differential equation that can be solved as follows:

$$\mathcal{T}'(t) = -\mathcal{T}(t) \cdot \sigma(t) \quad (3)$$

$$\frac{\mathcal{T}'(t)}{\mathcal{T}(t)} = -\sigma(t) \quad (4)$$

$$\int_a^b \frac{\mathcal{T}'(t)}{\mathcal{T}(t)} dt = - \int_a^b \sigma(t) dt \quad (5)$$

$$\log \mathcal{T}(t)|_a^b = - \int_a^b \sigma(t) dt \quad (6)$$

$$\mathcal{T}(a \rightarrow b) \equiv \frac{\mathcal{T}(b)}{\mathcal{T}(a)} = \exp \left(- \int_a^b \sigma(t) dt \right) \quad (7)$$

where we define $\mathcal{T}(a \rightarrow b)$ as the probability that the ray travels from distance a to b along the ray without hitting a particle, which is related to the previous notation by $\mathcal{T}(t) = \mathcal{T}(0 \rightarrow t)$.

¹If you are interested in borrowing the LaTeX notation, please refer to: <https://www.overleaf.com/read/fkxpkzxhnyws>

Probabilistic interpretation. Note that we can also interpret the function $1-\mathcal{T}(t)$ (often called *opacity*) as a cumulative distribution function (CDF) indicating the probability that the ray *does* hit a particle sometime before reaching distance t . Then $\mathcal{T}(t) \cdot \sigma(t)$ is the corresponding probability density function (PDF), giving the likelihood that the ray stops precisely at distance t .

Volume rendering. We can now calculate the expected value of the light emitted by the particles in the volume as the ray travels from $t=0$ to D , composited on top of a background color. Since the probability density for stopping at t is $\mathcal{T}(t) \cdot \sigma(t)$, the expected color is

$$\mathbf{C} = \int_0^D \mathcal{T}(t) \cdot \sigma(t) \cdot \mathbf{c}(t) dt + \mathcal{T}(D) \cdot \mathbf{c}_{\text{bg}} \quad (8)$$

where \mathbf{c}_{bg} is a background color that is composited with the foreground scene according to the residual transmittance $\mathcal{T}(D)$. Without loss of generality, we omit the background term in what follows.

Homogeneous media. We can calculate the color of some homogeneous volumetric media with constant color \mathbf{c}_a and density σ_a over a ray segment $[a, b]$ by integration:

$$\mathbf{C}(a \rightarrow b) = \int_a^b \mathcal{T}(a \rightarrow t) \cdot \sigma(t) \cdot \mathbf{c}(t) dt \quad (9)$$

$$= \sigma_a \cdot \mathbf{c}_a \int_a^b \mathcal{T}(a \rightarrow t) dt \quad \text{constant density/radiance} \quad (10)$$

$$= \sigma_a \cdot \mathbf{c}_a \int_a^b \exp\left(-\int_a^t \sigma(u) du\right) dt \quad \text{substituting (7)} \quad (11)$$

$$= \sigma_a \cdot \mathbf{c}_a \int_a^b \exp\left(-\sigma_a u \Big|_a^t\right) dt \quad \text{constant density (again)} \quad (12)$$

$$= \sigma_a \cdot \mathbf{c}_a \int_a^b \exp(-\sigma_a(t-a)) dt \quad (13)$$

$$= \sigma_a \cdot \mathbf{c}_a \cdot \frac{\exp(-\sigma_a(t-a))}{-\sigma_a} \Big|_a^b \quad (14)$$

$$= \mathbf{c}_a \cdot (1 - \exp(-\sigma_a(b-a))) \quad (15)$$

Transmittance is multiplicative. Note that transmittance factorizes as follows:

$$\mathcal{T}(a \rightarrow c) = \exp\left(-\left[\int_a^b \sigma(t) dt + \int_b^c \sigma(t) dt\right]\right) \quad (16)$$

$$= \exp\left(-\int_a^b \sigma(t) dt\right) \exp\left(-\int_b^c \sigma(t) dt\right) \quad (17)$$

$$= \mathcal{T}(a \rightarrow b) \cdot \mathcal{T}(b \rightarrow c) \quad (18)$$

This also follows from the probabilistic interpretation of \mathcal{T} , since the probability that the ray does not hit any particles within $[a, c]$ is the product of the probabilities of the two independent events that it does not hit any particles within $[a, b]$ or within $[b, c]$.

Transmittance for piecewise constant data. Given a set of intervals $\{[t_n, t_{n+1}]\}_{n=1}^N$ with constant density σ_n within the n -th segment, and with $t_1=0$ and $\delta_n=t_{n+1}-t_n$, transmittance is equal to:

$$\mathcal{T}_n = \mathcal{T}(t_n) = \mathcal{T}(0 \rightarrow t_n) = \exp\left(-\int_0^{t_n} \sigma(t) dt\right) = \exp\left(\sum_{k=1}^{n-1} -\sigma_k \delta_k\right) \quad (19)$$

Volume rendering of piecewise constant data. Combining the above, we can evaluate the volume rendering integral through a medium with piecewise constant color and density:

$$\mathbf{C}(t_{N+1}) = \sum_{n=1}^N \int_{t_n}^{t_{n+1}} \mathcal{T}(t) \cdot \sigma_n \cdot \mathbf{c}_n dt \quad \text{piecewise constant} \quad (20)$$

$$= \sum_{n=1}^N \int_{t_n}^{t_{n+1}} \mathcal{T}(0 \rightarrow t_n) \cdot \mathcal{T}(t_n \rightarrow t) \cdot \sigma_n \cdot \mathbf{c}_n dt \quad \text{from (18)} \quad (21)$$

$$= \sum_{n=1}^N \mathcal{T}(0 \rightarrow t_n) \int_{t_n}^{t_{n+1}} \mathcal{T}(t_n \rightarrow t) \cdot \sigma_n \cdot \mathbf{c}_n dt \quad \text{constant} \quad (22)$$

$$= \sum_{n=1}^N \mathcal{T}(0 \rightarrow t_n) \cdot (1 - \exp(-\sigma_n(t_{n+1} - t_n))) \cdot \mathbf{c}_n \quad \text{from (15)} \quad (23)$$

This leads to the volume rendering equations from NeRF [3, Eq.3]:

$$\mathbf{C}(t_{N+1}) = \sum_{n=1}^N \mathcal{T}_n \cdot (1 - \exp(-\sigma_n \delta_n)) \cdot \mathbf{c}_n, \quad \text{where} \quad \mathcal{T}_n = \exp\left(\sum_{k=1}^{n-1} -\sigma_k \delta_k\right) \quad (24)$$

Finally, rather than writing these expressions in terms of volumetric density, we can re-express them in terms of alpha-compositing weights $\alpha_n \equiv 1 - \exp(-\sigma_n \delta_n)$, and by noting that $\prod_i \exp x_i = \exp(\sum_i x_i)$ in (19):

$$\mathbf{C}(t_{N+1}) = \sum_{n=1}^N \mathcal{T}_n \cdot \alpha_n \cdot \mathbf{c}_n, \quad \text{where} \quad \mathcal{T}_n = \prod_{i=1}^{n-1} (1 - \alpha_i) \quad (25)$$

Alternate derivation. By making use of the earlier connection between CDF and PDF that $(1 - \mathcal{T})' = \mathcal{T} \sigma$, and by assuming constant color \mathbf{c}_a along an interval $[a, b]$:

$$\int_a^b \mathcal{T}(t) \cdot \sigma(t) \cdot \mathbf{c}(t) dt = \mathbf{c}_a \int_a^b (1 - \mathcal{T})'(t) dt \quad (26)$$

$$= \mathbf{c}_a \cdot (1 - \mathcal{T}(t))|_a^b \quad (27)$$

$$= \mathbf{c}_a \cdot (\mathcal{T}(a) - \mathcal{T}(b)) \quad (28)$$

$$= \mathbf{c}_a \cdot \mathcal{T}(a) \cdot (1 - \mathcal{T}(a \rightarrow b)) \quad (29)$$

Combined with constant per-interval density, this identity yields the same expression for color as (24).

References

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- [4] Matt Pharr, Wenzel Jakob, and Greg Humphreys. *Physically based rendering: From theory to implementation*. Morgan Kaufmann, 2016. https://pbr-book.org/3ed-2018/Volume_Scattering/Volume_Scattering_Processes.

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